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Work supervised by F. G. Scholtz

ROOM 316 TALK:



PINHOLE INTERFERENCE IN 3D FUZZY SPACE

A natural quantum-to-classical transition

Stellenbosch University
August 2023

Outline

1 SETUP

- Motivation & goal
- Existing work

2 Fuzzy space FORMALISM

- States & observables
- Position measurement
- Coordinate representation

3 FREE PARTICLE solutions

- Plane & spherical waves

4 INTERFERENCE calculation

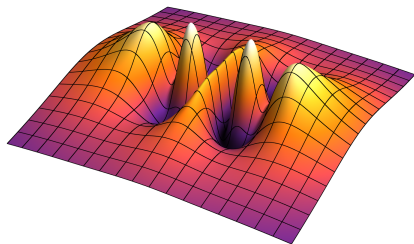
- Commutative & non-commutative

5 DISCUSSION

- Interference suppression
- Many-particles

6 SUMMARY

THE SETUP





Question:

Can the structure of spacetime at **SMALLEST LENGTH SCALE** affect the physics we perceive at **LARGER LENGTH SCALES** (*i.e. classical physics*)?



Our investigation:

Study **DOUBLE-PINHOLE SETUP** in the 3D **FUZZY SPACE** formalism of **NON-COMMUTATIVE** quantum mechanics.



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- **MINIMUM LENGTH** scale \longleftarrow set by a parameter λ
- Otherwise **ORDINARY QUANTUM MECHANICS !**

Motivation & Goal

Why this particular investigation?



Why non-commutative geometry?



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- Doplicher et al. (1995):

Non-trivial **SMALL-SCALE** spacetime structure (esp. min length)
likely necessary for **QM + GRAVITY**



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- Illustrative **TOY MODEL** ← quantum behaviour = interference
- **QUANTIFIABLE SUPPRESSION** strength
- Good setup for **EXPERIMENTAL TESTING**



Pittaway & Scholtz (2021)

2-SLITS in non-commutative 2D MOYAL PLANE



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Interference suppression at:

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- large PARTICLE NUMBER, N



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- Moyal plane interference pattern:
 - is **ASYMMETRIC** under reflection ← ∴ Moyal commutation relations break rotational symmetry
 - has **UNOBSERVABLE SUPPRESSION**



Moyal plane interference

$$P(\mathbf{D}) = 1 + \underbrace{\exp\left[-\frac{N\theta m^2}{2\hbar^2} \mathbf{v}^2 (1 - \cos \alpha)\right]}_{\text{interference suppression}} \underbrace{\cos(\dots)}_{\text{interference}}$$



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- $\sqrt{\theta} \sim 10^{-35}$ m ← Planck length
- $N \sim 10^{23}$ ← Avogadro's number



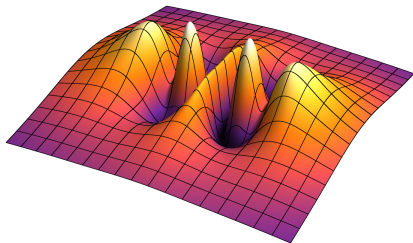
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THE FORMALISM





Definition (Fuzzy space)

$$\mathbf{1} \text{ } \mathfrak{su}(2) \text{ ALGEBRA} \longrightarrow [\hat{x}_i, \hat{x}_j] = 2i\lambda \epsilon_{ijk} \hat{x}_k$$



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i.e. $\hat{r}|n_1, n_2\rangle = \lambda(n_1 + n_2 + 1)|n_1, n_2\rangle$



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\implies 1 copy of each quantised radius



Definition (Quantum state space)

1 HILBERT SPACE

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so $\left\| \text{Proj}_{[1] \oplus \dots \oplus [M]} \right\|$
 $\sim \frac{4}{3}\pi [\lambda(N+1)]^3$



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HERMITIAN OPERATORS on \mathcal{H}_q



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Important observables

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■ HAMILTONIAN $\longrightarrow \hat{H} = -\frac{\hbar^2}{2m} \hat{\Delta} + V(\hat{R}) \longleftarrow$ like normal QM



Physical subspace

$$\mathcal{H}_q = \bigoplus_{n \in \mathbb{N}} [n] \otimes [n]^* \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$$

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3 PROJECTION $\longrightarrow \hat{Q} := \frac{1}{2\pi} \int_0^{2\pi} e^{i\phi \hat{\Gamma}} d\phi$



Definition (Position measurement)

1 Position **EIGENSTATES** \longrightarrow **MINIMUM-UNCERTAINTY** states

$$|\mathbf{z}\rangle := e^{-\frac{1}{2}\bar{z}_\alpha z_\alpha} e^{z_\alpha a_\alpha^\dagger} |0\rangle, \text{ for}$$
$$\mathbf{z} := e^{i\gamma} \sqrt{\frac{r}{\lambda}} \begin{bmatrix} \cos\left(\frac{\theta}{2}\right) e^{-i\frac{\phi}{2}} \\ \sin\left(\frac{\theta}{2}\right) e^{i\frac{\phi}{2}} \end{bmatrix} \in \mathbb{C}^2 \longleftarrow \text{encodes } \mathbf{D} = (r, \theta, \phi)$$



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2 **POVM**

$$\begin{aligned}
 \rightarrow \quad |z_1, z_2, n_1, n_2\rangle_{\text{ph}} &:= \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2\hat{r}}} |\mathbf{z}\rangle \langle n_1, n_2| \\
 \rightarrow \quad \hat{\Pi}_{\mathbf{z}} &:= \sum_{n_1, n_2} |z_1, z_2, n_1, n_2\rangle_{\text{ph}} \langle n_1, n_2|
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3 **BORN RULE** $\longrightarrow P(\mathbf{D}) = \text{Tr}_q(\hat{\Pi}_{\mathbf{z}}\rho)$



Weak measurement picture

$$\mathbf{1} \quad \mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$$



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- 1 $\mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$
- 2 \hat{X}_i act only on \mathcal{H}_c



Weak measurement picture

- 1 $\mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$
- 2 \hat{X}_i act only on \mathcal{H}_c
- 3 \therefore **POSITION** measurement
 - **LOCAL** measurement
 - traces out “**ENVIRONMENT**”, \mathcal{H}_c^*



Definition (Symbol & star product)

1 POSITION-ENCODING states $\rightarrow |\mathbf{z}\rangle := \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2\hat{r}}} |\mathbf{z}\rangle\langle\mathbf{z}|$



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ψ is indep of γ
 \therefore function on \mathbb{R}^3



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3 **SYMBOL** $\longrightarrow \langle\mathbf{z}|\psi|\mathbf{z}\rangle$



Definition (Symbol & star product)

1 **POSITION-ENCODING** states $\longrightarrow |\mathbf{z}\rangle := \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2\hat{r}}} |\mathbf{z}\rangle\langle\mathbf{z}|$

2 **COORDINATE REP** $\longrightarrow \psi(\mathbf{z}) := (\mathbf{z}|\psi) = \langle\mathbf{z}|\sqrt{4\pi\lambda^2\hat{r}}\psi|\mathbf{z}\rangle$

3 **SYMBOL** $\longrightarrow \langle\mathbf{z}|\psi|\mathbf{z}\rangle$

4 **VOROS PRODUCT** $\longrightarrow \star := \exp\left[\overleftarrow{\partial}_{z_\alpha} \overrightarrow{\partial}_{\bar{z}_\alpha}\right]$



Notable properties

1 COMPLETENESS $\longrightarrow \int \frac{d^4 z}{\pi^2} |\mathbf{z}\rangle \bar{\mathbf{x}}(\mathbf{z}| = \mathbf{1}_q$



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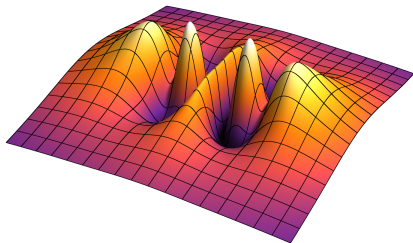
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\implies Alternate “wave-mechanics” development !



FREE PARTICLE SOLUTIONS





Non-commutative free Schrödinger equation

$$\hat{H}|\psi\rangle = -\frac{\hbar^2}{2m}\hat{\Delta}|\psi\rangle = E|\psi\rangle$$



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Types of solutions:

1 PLANE WAVE $\longrightarrow |\mathbf{k}\rangle := e^{i\mathbf{k}\cdot\hat{\mathbf{x}}}$ \longleftarrow typical form

2 SPHERICAL WAVE $\longrightarrow |k, l, m\rangle = g_l(\hat{n} - l, k) \hat{Y}_{lm}$



Plane wave solutions

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$$\cos(\lambda k_3) = \cos(\lambda k_1) \cos(\lambda k_2) - \hat{\mathbf{k}}_1 \cdot \hat{\mathbf{k}}_2 \sin(\lambda k_1) \sin(\lambda k_2),$$

$$\sin(\lambda k_3) \hat{\mathbf{k}}_3 = \hat{\mathbf{k}}_1 \sin(\lambda k_1) \cos(\lambda k_2) + \hat{\mathbf{k}}_2 \sin(\lambda k_2) \cos(\lambda k_1) - \hat{\mathbf{k}}_1 \times \hat{\mathbf{k}}_2 \sin(\lambda k_1) \sin(\lambda k_2)$$

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- DISPERSION RELATION $\longrightarrow \hat{H} |\mathbf{k}\rangle = \frac{2\hbar^2}{m\lambda^2} \sin^2\left(\frac{k\lambda}{2}\right) |\mathbf{k}\rangle$

Spherical wave solutions



2 SPHERICAL WAVE $\longrightarrow |k, l, m\rangle = g_l(\hat{n} - l, k) \hat{Y}_{lm}$



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Properties

■ Permit 2 RADIAL SOLUTIONS $\longrightarrow g_l = A g_{J,l} + B g_{Y,l}$
c.f. spherical Bessel- & Neumann



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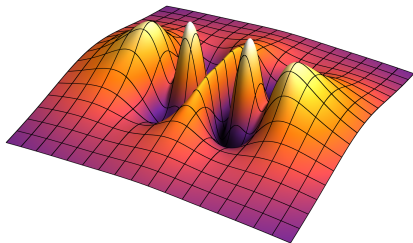
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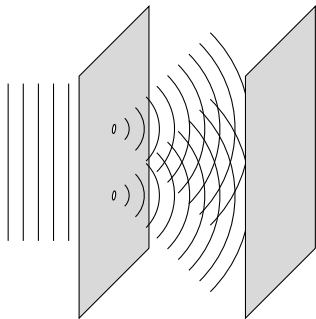
$$\therefore \langle \mathbf{z} | g_{H,l}(\hat{n}, k) | \mathbf{z} \rangle \sim \frac{e^{r(\cos(\lambda k) + i \sin(\lambda k) - 1)/\lambda}}{(ir/\lambda)^{l+1}}$$

MAIN CALCULATION



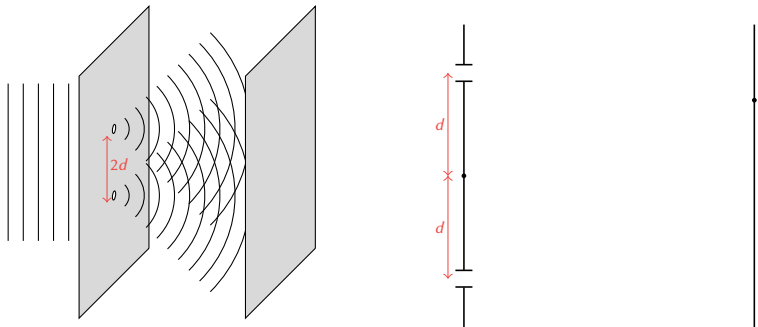
Double-Pinhole Setup

The most famous quantum experiment



Double-Pinhole Setup

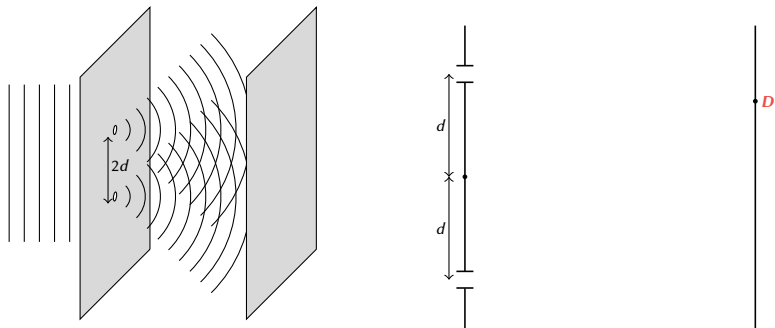
The most famous quantum experiment



■ PINHOLES $\longrightarrow z = \pm d$

Double-Pinhole Setup

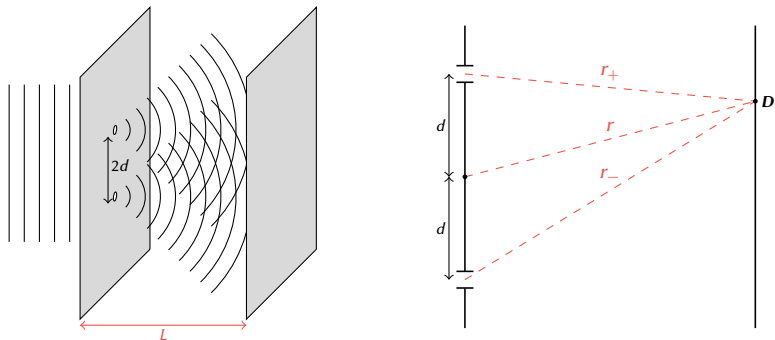
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- **PINHOLES** $\longrightarrow z = \pm d$
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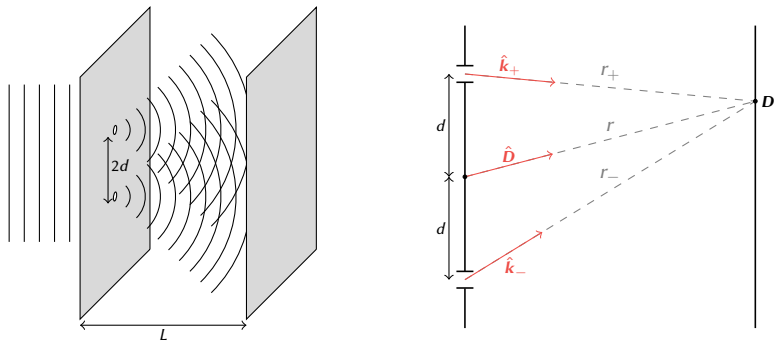
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- **PINHOLES** $\longrightarrow z = \pm d$
- **DETECTION POINT** $\longrightarrow \mathbf{D} = (L, y_D, z_D) \equiv (r, \theta, \phi)$
- **DISTANCES** $\longrightarrow L, r, r_{\pm}$ \longleftarrow each $\gg d$: large separation approx

Double-Pinhole Setup

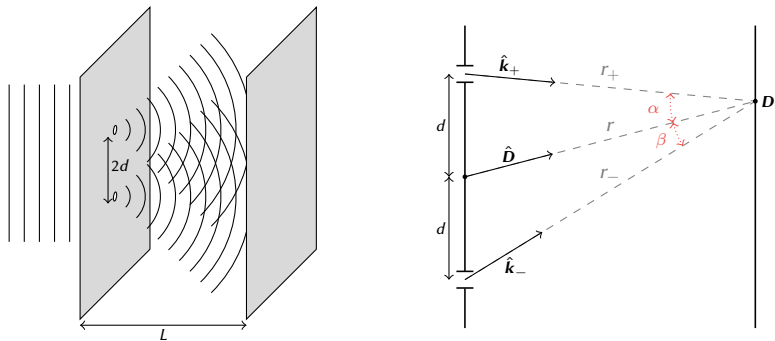
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- **ANGLES** $\longrightarrow \alpha, \beta$



Interference calculation overview

$$\mathbf{1} \text{ STATE at } \mathbf{D} \longrightarrow \psi \sim \left\{ \begin{array}{l} \text{spherical wave} \\ \text{from } z = +d \end{array} \right\} + \left\{ \begin{array}{l} \text{spherical wave} \\ \text{from } z = -d \end{array} \right\}$$



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Commutative interference calculation

1 STATE at \mathbf{D} \longrightarrow $\psi(\mathbf{D}) \sim \frac{1}{r_+} e^{ikr_+} + \frac{1}{r_-} e^{ikr_-}$ \longleftarrow asymptotic form of spherical Hankel



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$$P_{\text{comm}}(\mathbf{D}) \sim \frac{1}{r_+ r_-} \left[\underbrace{\frac{2d^2}{r_+ r_-} + \cos(\alpha + \beta)}_{\text{bimodal shaping function}} + \underbrace{\cos(rk(\cos \alpha - \cos \beta))}_{\text{interference terms}} \right]$$



Non-commutative interference calculation

$$\mathbf{1} \text{ SYMBOL at } \mathbf{D} \longrightarrow \langle \mathbf{z} | \psi | \mathbf{z} \rangle \sim \langle \mathbf{z}^+ | g_k(\hat{n}) | \mathbf{z}^+ \rangle + \langle \mathbf{z}^- | g_k(\hat{n}) | \mathbf{z}^- \rangle$$



Non-commutative interference calculation

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- 2 **PARAXIAL** approximation \longrightarrow $\langle \mathbf{z}^\pm | g_k(\hat{n}) | \mathbf{z}^\pm \rangle \sim \langle \mathbf{z} | \eta_\pm e^{ik^\pm \cdot \hat{\mathbf{x}}} | \mathbf{z} \rangle$



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- 4 **BORN RULE** $\longrightarrow P(\mathbf{D}) = \text{Tr}_q \left(\hat{\Pi}_{\mathbf{z}} |\psi\rangle\langle\psi| \right)$
- 5 Compute remaining **MATRIX ELEMENTS**

Interference Calculation

The main result!



$$P(\mathbf{D}) \sim \underbrace{\frac{\eta_+^2 + \eta_-^2}{2} \left(\frac{r}{\lambda} + 1 \right)}_{\text{bimodal shaping function}} + \underbrace{\eta_+ \eta_- e^{A-r/\lambda}}_{\text{interference suppression}} \times \underbrace{((A+1) \cos B - B \sin B)}_{\text{interference terms}}$$

$$\text{where } \begin{cases} \eta_{\pm} & := \frac{\lambda}{r_{\pm}} \exp\left[\frac{1}{\lambda}(r_{\pm} - r)(\cos(\lambda k) - 1)\right], \\ A & := \frac{r}{\lambda} (\cos^2(\lambda k) + \cos(\alpha + \beta) \sin^2(\lambda k)), \\ B & := \frac{r}{\lambda} \sin(\lambda k) \cos(\lambda k) (\cos \alpha - \cos \beta) \end{cases}$$

Interference Calculation

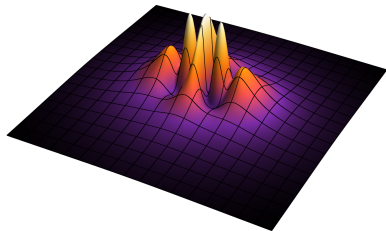
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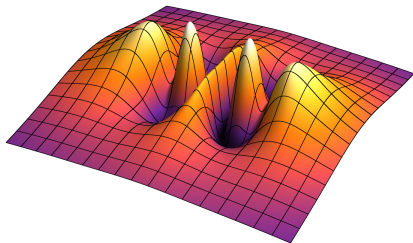
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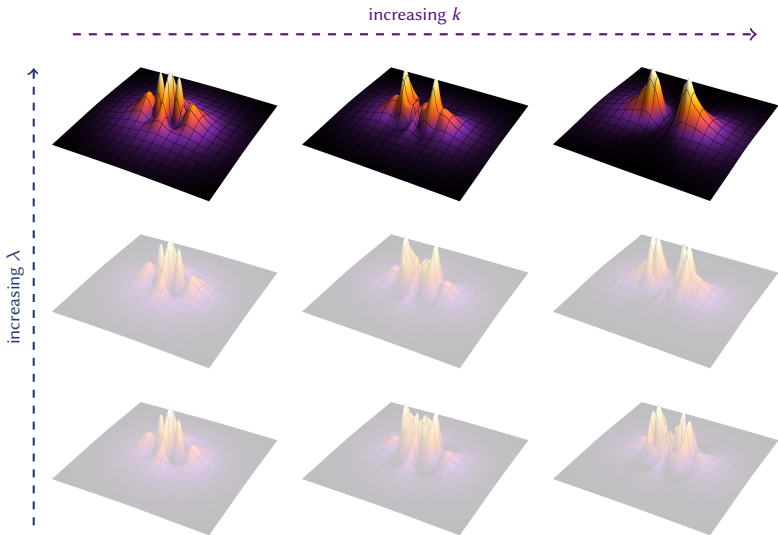


DISCUSSION OF RESULTS



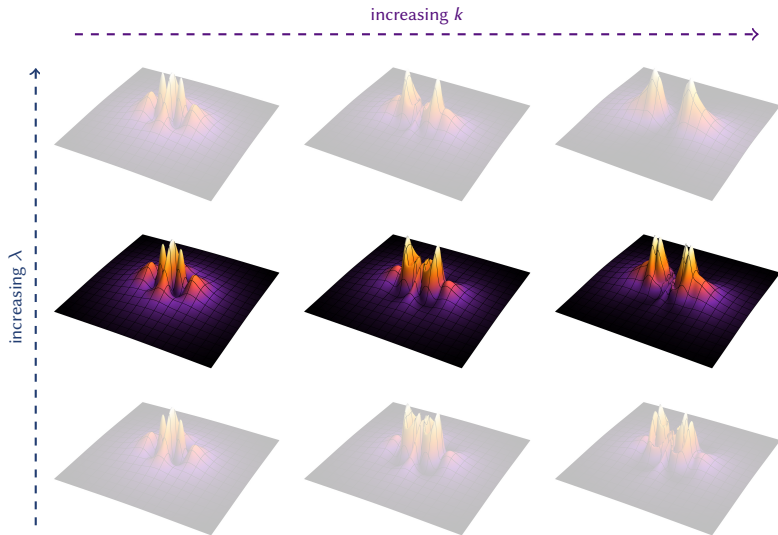
Qualitative Behaviour

& commutative limit



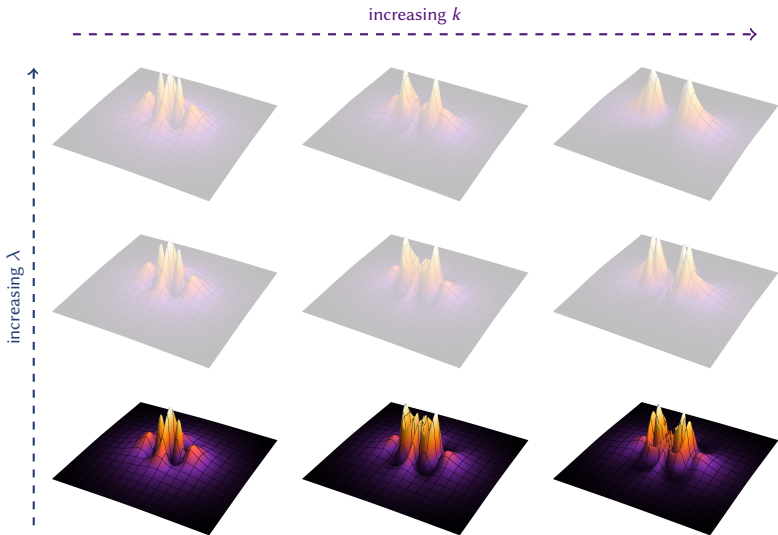
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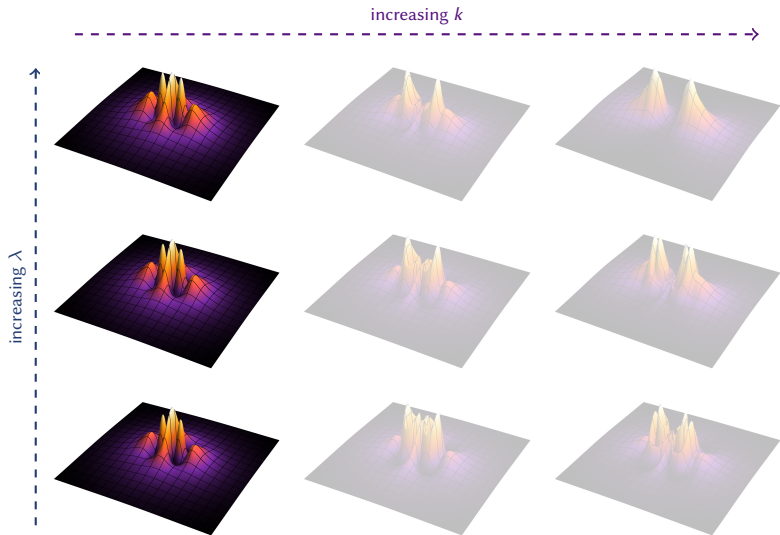
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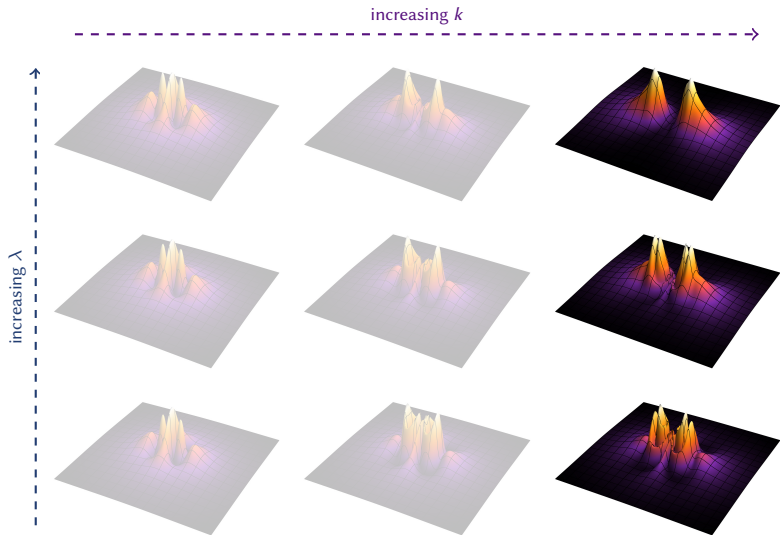
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Fuzzy space interference

$$P(\mathbf{D}) \sim \dots + \underbrace{\eta_+ \eta_- \exp\left[\frac{r}{\lambda} \sin^2(\lambda k)(\cos(\alpha + \beta) - 1)\right]}_{\text{interference suppression}} \underbrace{(\dots)}_{\text{interference}}$$



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- $m \sim 10^{-31} \text{ kg}$ ← mass of electron
- $\lambda \sim 10^{-35} \text{ m}$ ← Planck length



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- } $\implies r \lesssim 10^{-21} \text{ m} !$



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Important features

- r dependence
- suppression possible at low k



N particles

1 HILBERT SPACE $\longrightarrow \mathcal{H}_q^{\text{tot}} = \bigotimes_{n=1}^N \mathcal{H}_q^{(n)}$



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- 3 Free HAMILTONIAN $\longrightarrow \hat{H}_{\text{free}}^{\text{tot}} = -\frac{\hbar^2}{2m} \sum_{n=1}^N \hat{\Delta}^{(n)}$



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1 HILBERT SPACE $\longrightarrow \mathcal{H}_q^{\text{tot}} = \bigotimes_{n=1}^N \mathcal{H}_q^{(n)}$

2 ALGEBRA $\longrightarrow [\hat{x}_i^{(l)}, \hat{x}_j^{(n)}] = 2i\lambda \epsilon_{ijk} \delta_{ln} \hat{x}_k^{(l)}$

3 Free HAMILTONIAN $\longrightarrow \hat{H}_{\text{free}}^{\text{tot}} = -\frac{\hbar^2}{2m} \sum_{n=1}^N \hat{\Delta}^{(n)}$

4 PLANE WAVES $\longrightarrow |\mathbf{k}^{(i \dots N)}\rangle = \exp \left[i \sum_{n=1}^N \mathbf{k}^{(n)} \cdot \hat{\mathbf{x}}^{(n)} \right]$

\vdots



N particles

5 CENTER-OF-MASS frame

$$\longrightarrow \hat{\mathbf{x}}^{(\text{CM})} := \frac{1}{N} \sum_{n=1}^N \hat{\mathbf{x}}^{(n)}, \quad \hat{\boldsymbol{\xi}}^{(n)} := \hat{\mathbf{x}}^{(n)} - \hat{\mathbf{x}}^{(\text{CM})}$$

$$\mathbf{k}^{\text{tot}} := \sum_{n=1}^N \mathbf{k}^{(n)}, \quad \mathbf{q}^{(n)} := \mathbf{k}^{(n)} - \frac{1}{N} \mathbf{k}^{\text{tot}}$$

⋮



N particles

6 SPLIT Hamiltonian* \longrightarrow \hat{H}^{tot}



N particles

6 SPLIT Hamiltonian*

→

\hat{H}^{tot}

$\hat{H}_{\text{free}}^{\text{tot}}$

+

$\hat{H}_{\text{interaction}}^{\text{tot}}$



N particles

6 **SPLIT** Hamiltonian* \rightarrow \hat{H}^{tot}

$$\hat{H}_{\text{free}}^{\text{tot}} + \hat{H}_{\text{interaction}}^{\text{tot}} = 0$$



N particles

6 **SPLIT** Hamiltonian* \rightarrow

$$\hat{H}_{\text{CoM}} + \hat{H}_{\text{internal}} = \hat{H}_{\text{free}}^{\text{tot}} + \hat{H}_{\text{interaction}}^{\text{tot}} = 0$$

The diagram illustrates the decomposition of the total Hamiltonian \hat{H}^{tot} for N particles. It shows a hierarchical structure where the total Hamiltonian is split into a free part and an interaction part, which are further decomposed into center-of-mass and internal components. Dashed red lines connect the terms to show their relationships.



N particles

6 **SPLIT** Hamiltonian* \rightarrow \hat{H}^{tot}

$$\hat{H}_{\text{CoM}} + \hat{H}_{\text{internal}} \xrightarrow{\text{dashed lines}} \hat{H}_{\text{free}}^{\text{tot}} + \hat{H}_{\text{interaction}}^{\text{tot}} = 0$$

$\mathbf{q}^{(n)} = 0$



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$$\Rightarrow \text{CoM dynamics} \leftrightarrow \text{1 PARTICLE: } \left\{ \begin{array}{ll} \text{mass} & Nm \\ \text{momentum} & \mathbf{k}^{\text{tot}} \\ \text{NC param} & \lambda/N \end{array} \right\}$$



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7 **NEGLECT** corrections \rightarrow **EXPAND** in λ and $T \sim \frac{\hbar^2 (q^{(n)})^2}{mk_B}$



Fuzzy space interference

$$P(\mathbf{D}) \sim \dots + \eta_+ \eta_- \exp \left[\underbrace{\frac{Nr}{\lambda} \sin^2(\lambda k^{\text{tot}}/N) (\cos(\alpha + \beta) - 1)}_{\text{interference suppression}} \right] \underbrace{(\dots)}_{\text{interference}}$$



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- $\langle E \rangle \sim 1 \text{ eV}$
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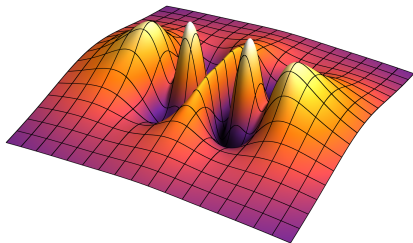
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CONCLUDING REMARKS





Challenges

- 1 **CREATE** \otimes **MANIPULATE** massive quantum superposition

Experimental Prospects

How might we observe suppression practically?



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- 2 ISOLATE from TRUE ENVIRONMENT

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 - *van Es et al. (2008)* → split propagating 10^4 -atom ^{87}Rb BEC
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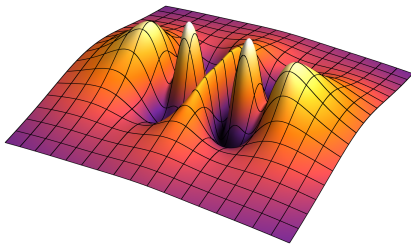
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- 3 Macroscopic OPTOMECHANICAL superposition ?
 - *Kleckner et al. (2008)* → describe entangling photon with macroscopic cantilever



Key takeaways

- Fuzzy space \longrightarrow **CLASSICAL TRANSITION** without heat bath
- Quantum suppression \longrightarrow realistically **OBSERVABLE**
- Suppression strength \longrightarrow **EXTENSIVE & DISTANCE**-dependent





1 ALTERNATE SETUPS:

- Treat fuzzy-space **VON NEUMANN MEASUREMENT**

2 FORMALISM EXTENSIONS:

- Extend to non-commutative **QFT**

3 EXPERIMENTAL VERIFICATION:

- Implement proposed experiment
- Devise alternate experiment



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