Exploring **Tensor Products**



"A tensor is a thing that transforms like a tensor!"

Stellenbosch University May 2021 20854714



Motivation & Overview

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Where we are going...

Why discuss tensors?

Dario Trinchero Exploring tensor products





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commonly encountered in undergrad without careful definition





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Where we are heading

We build to defining the tensor product as the unique space linearizing bilinear maps. We touch on more abstract settings, and the tensor-hom adjunction.





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•
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 \exists x^* \in X: x + x^* = 0_X$$
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$\bullet (\lambda + \mu) \cdot \mathbf{x} = \lambda \cdot \mathbf{x} + \mu \cdot \mathbf{x}$) distributive laws
$\bullet (\lambda \mu) \cdot \mathbf{x} = \lambda \cdot (\mu \cdot \mathbf{x})$	compatibility of \cdot with $\mathbb C$
$\bullet 1 \cdot x = x$	identity of \cdot





A complex VECTOR SPACE is an Abelian group (X, +) with operation $\therefore \mathbb{C} \times X \to X$, satisfying, for all $x, y \in X, \lambda, \mu \in \mathbb{C}$,

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Notation

Hereafter, take *X*, *Y*, *Z* to be complex vector spaces.





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Notation

Hereafter, take *X*, *Y*, *Z* to be complex vector spaces.

Definition (Dual space)

The vector space X^* of all linear functionals $X \to \mathbb{C}$ is the DUAL SPACE of X. If X has basis $\mathcal{B} = (b_1, \ldots, b_n)$, then X^* has DUAL BASIS $(\beta^1, \ldots, \beta^n)$, where $\beta^j(b_i) := \delta_{ij}$.



Preliminaries Equivalence relations & quotients



Definition (Equivalence relation)

An EQUIVALENCE RELATION \sim on a set *S* is a binary relation (subset of $S \times S$) satisfying, for all $x, y, z \in S$,





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 $x \sim x$ reflexivity $x \sim y \implies y \sim x$ symmetry $x \sim y \wedge y \sim z \implies x \sim z$ transitivity





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- $x \sim y \land y \sim z \implies x \sim z$ transitivity

Definition (Quotient vector space)

For each $x \in S$, define the EQUIVALENCE CLASS $[x] \coloneqq \{y \in S : y \sim x\}$, and the QUOTIENT SPACE $S / \sim \coloneqq \{[x] : x \in S\}$. For vector space X, note that X / \sim is also a vector space if we define $[x] + [y] \coloneqq [x + y], \lambda[x] \coloneqq [\lambda x]$.





Texts aimed at physicists[7] typically define:

Definition 1 (Tensor)

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$$[T_{\mathcal{C}}]_{j'_{1}\dots j'_{q}}^{i'_{1}\dots i'_{p}} = (P^{-1})_{i_{1}}^{i'_{1}}\cdots (P^{-1})_{i_{p}}^{i'_{p}} [T_{\mathcal{B}}]_{j_{1}\dots j_{q}}^{i_{1}\dots i_{p}} P_{j'_{1}}^{j_{1}}\cdots P_{j'_{q}}^{j_{q}}$$





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..... what??!





Example (Tensor of type (1, 0)) 1D array $[T_{\mathcal{B}}]^{i}$, which under $\mathcal{B} = (b_{1}, \ldots, b_{n}) \mapsto (b_{i}P_{1}^{i}, \ldots, b_{i}P_{n}^{i}) = \mathcal{C}$, transforms by $[T_{\mathcal{C}}]^{i'} = (P^{-1})_{i}^{i'} [T_{\mathcal{B}}]^{i}$





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recall change-of-basis matrix

$$P_{\mathcal{B}\leftarrow\mathcal{C}} := \begin{bmatrix} | & | & | \\ [c_1]_{\mathcal{B}} & [c_2]_{\mathcal{B}} & \dots & [c_n]_{\mathcal{B}} \\ | & | & | & | \end{bmatrix}$$





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• so the basis vectors transform $c_j = b_i (P_{\mathcal{B}\leftarrow \mathcal{C}})_{ij} \equiv b_i P_j^i$ from above





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• but coordinates transform by the inverse matrix: $[v_{\mathcal{C}}] = P_{\mathcal{C} \leftarrow \mathcal{B}} [v_{\mathcal{B}}] = (P_{\mathcal{B} \leftarrow \mathcal{C}})^{-1} [v_{\mathcal{B}}],$





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- the basis vectors transform $c_j = b_i (P_{\mathcal{B} \leftarrow \mathcal{C}})_{ij} \equiv b_i P_j^i$

• hence
$$[f_{\mathcal{C}}]_j = f(c_j) = f(b_i P_j^i) = f(b_i) P_j^i = [f_{\mathcal{B}}]_i P_j^i$$





Example (Tensor of type (1, 1)) 2D array $[T_{\mathcal{B}}]_{j}^{i}$, which under $\mathcal{B} = (b_{1}, \ldots, b_{n}) \mapsto (b_{i}P_{1}^{i}, \ldots, b_{i}P_{n}^{i}) = \mathcal{C}$, transforms by $[T_{\mathcal{C}}]_{j'}^{i'} = (P^{-1})_{i}^{i'} [T_{\mathcal{B}}]_{j}^{i}P_{j'}^{j}$





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■ the components [M_B]_{ij} of linear operator M with respect to basis B are computed [M_B]_{ij} = [M(b_j)_B]ⁱ





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Just the transformation of linear operator components (matrices)!

- the components [M_B]_{ij} of linear operator M with respect to basis B are computed [M_B]_{ij} = [M(b_j)_B]ⁱ
- applying previous examples, with linearity,

$$[M_{\mathcal{C}}]_{ij} = [M(c_j)_{\mathcal{C}}]^i = \left[\left(T(b_{j'}) P_j^{j'} \right)_{\mathcal{C}} \right]$$

= $(P^{-1})^i_{i'} [M(b_{j'})_{\mathcal{B}}]^{i'} P_j^{j'} = (P^{-1})^i_{i'} [M_{\mathcal{B}}]_{i'j'} P_j^{j'}$





Our examples suggest, informally,

- **1** contravariant (upper) indices act like vector components
- 2 covariant (lower) indices act like covector (linear functional) components





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Question: Type-(1, 1) tensors act like matrices. Can we view a matrix in terms of vector-covector pairs? Yes. Consider theorem 7.14 from Poole[5]:

Theorem (Singular-value decomposition) An $m \times n$ matrix A of rank $r \le n$ may be written in the form

$$A = \sigma_1 \mathbf{u}_1 \mathbf{v}_1^T + \dots + \sigma_r \mathbf{u}_r \mathbf{v}_r^T,$$

where the σ_i are called the "singular values" of A.



The Free Vector Space Brief interlude to set up tensor product



We can construct a vector space with any given set as basis:

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Note (Formal linear combinations)

We commonly write $f \in F(S)$ as a formal linear combination

$$\sum_{x\in S}f(x)x$$

of elements of S. There are finitely-many non-0 coefficients, since f has finite support.





Definition 2 (Tensor product)

The TENSOR PRODUCT $X \otimes Y$ is the quotient space, $F(X \times Y) / \sim$, of the free vector space by the equivalence relation \sim generated by:

$$egin{aligned} & (x_1,y) + (x_2,y) \sim (x_1 + x_2,y), & (x,y_1) + (x,y_2) \sim (x,y_1 + y_2) \ & \lambda(x,y) \sim (\lambda x,y), & \lambda(x,y) \sim (x,\lambda y) \end{aligned}$$





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Write $x \otimes y := [(x, y)] \in X \otimes Y$.





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 $(\mathbf{x}, \mathbf{y}_1) + (\mathbf{x}, \mathbf{y}_2) \sim (\mathbf{x}, \mathbf{y}_1 + \mathbf{y}_2)$
 $\lambda(\mathbf{x}, \mathbf{y}) \sim (\lambda \mathbf{x}, \mathbf{y}),$ $\lambda(\mathbf{x}, \mathbf{y}) \sim (\mathbf{x}, \lambda \mathbf{y})$

Write $x \otimes y := [(x, y)] \in X \otimes Y$.

Definition 3 (Tensor) A TENSOR OF TYPE (p, q) on X is a vector in $\underbrace{X \otimes \cdots \otimes X}_{p} \otimes \underbrace{X^* \otimes \cdots \otimes X^*}_{q}$. This formalizes the link between (co-/)contravariant indices & (co-)vectors.





Lemma (Basis for tensor product)

Given bases $\mathcal{B} = (b_1, \ldots, b_n), \mathcal{C} = (c_1, \ldots, c_m)$ for X and Y respectively, the vectors $b_i \otimes c_j$ form a basis for $X \otimes Y$.

Proof.





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1 For linear indepdendence, note that $(b_i, c_j) \not\sim (b_{i'}, c_{j'})$ whenever $(i, j) \neq (i', j')$, or we would have two linearly-dependent vectors in either \mathcal{B} or \mathcal{C} .





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- **1** For linear indepdendence, note that $(b_i, c_j) \not\sim (b_{i'}, c_{j'})$ whenever $(i, j) \neq (i', j')$, or we would have two linearly-dependent vectors in either \mathcal{B} or \mathcal{C} .
- **2** The general element of $X \otimes Y$ has the form $\sum_i \lambda_i x_i \otimes y_i$, for $x_i \in X, y_i \in Y$. Expanding each x_i and y_i in the bases, and using linearity, we express this as a linear combination of vectors $b_i \otimes c_j$.





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Definition (Components of tensor)

The COMPONENTS $[T_{\mathcal{B}}]_{j_1...,j_q}^{l_1...,l_p}$ of tensor T of type (p, q) on X are its expansion coefficients with respect to the induced basis on $X \otimes \cdots \otimes X \otimes X^* \otimes \cdots \otimes X^*$:

$$T = [T_{\mathcal{B}}]_{j_1\ldots j_p}^{i_1\ldots i_p} b_{i_1}\otimes \cdots \otimes b_{i_p}\otimes \beta^{j_1}\otimes \cdots \otimes \beta^{j_q}.$$





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Note (Transformation law)

It is clear that this satisfies the required transformation law under change-of-basis. Loosely, each *b* vector acquires a prefactor of P^{-1} , and each β a prefactor of *P*.





Definition (Rank)

An ELEMENTARY TENSOR *R* (of type (p, q)) can be written as a single term $R = a_1 \otimes \cdots \otimes a_p \otimes \alpha_1 \otimes \cdots \otimes \alpha_q$, for non-0 $a_i \in X, \alpha_i \in X^*$. The RANK of tensor *T* is the minimum number of elementary tensors summing to *T*.





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Note (Rank)

This definition coincides with how the matrix rank (dimension of row-space) entered the singular value decomposition! A rank-*r* matrix was expressible thereby as a sum of *r* elementary terms $\sigma_i \mathbf{u}_i \mathbf{v}_i^T$.





Definition (Bilinear maps & forms)

A BILINEAR MAP is a map $F: X \times Y \rightarrow Z$, which is linear in both arguments:

$$F(\lambda x_1 + \mu x_2, y) = \lambda F(x_1, y) + \mu F(x_2, y)$$

$$F(x, \lambda y_1 + \mu y_2) = \lambda F(x, y_1) + \mu F(x, y_2)$$





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Write $B(X \times Y, Z)$ for the vector space of bilinear maps into Z, and $B(X \times Y)$ for that of bilinear forms.





Just as linear functionals are type-(0, 1) tensors, we can think of bilinear forms on $X \times X$ as type-(0, 2) tensors on X. This motivates the definition in Jeevanjee[1]:

Bilinear Form Definitions

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The definition in Ryan[6] looks subtly different:

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Note (Linearizing bilinear forms) Consider Ryan's definition of $x \otimes y \in B(X \times Y)^*$:



Linearizing Bilinear Maps



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Linking our Definitions



Note (Linking via quotient space)

Both definitions are manifestly spanned by pairs $x \otimes y$, and both produce pairs bilinear in $x \otimes y$. We can thus exhibit obvious isomorphisms from each of them to $F(X \times Y) / \sim$, proving all 3 definitions isomorphic.



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Alternatively, we can use the linearizing properties we showed. Use subsript R for Ryan and J for Jeevanjee.

Note (Linking via linearization)

The map $X \times Y \to B(X^* \times Y^*)$: $(x, y) \mapsto x \otimes_J y$, with $(x \otimes_J y)(f, g) := f(x)g(y)$, is manifestly bilinear. Hence there is unique linear map $l: X \otimes_R Y \to B(X^* \times Y^*)$, easily shown to be injective. So $X \otimes_R Y$ embeds in $B(X^* \times Y^*)$, & is isomorphic to the subspace spanned by $x \otimes_J y$.



More abstractly,

15/20

Definition 6 (Tensor product)

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Universal Property Definition Unifying all prior definitions

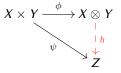
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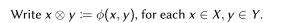
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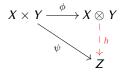
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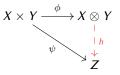


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Lemma (Uniqueness of $X \otimes Y$) $X \otimes Y$ is uniquely defined, up to isomorphism, by the above. Proof.

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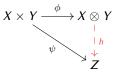
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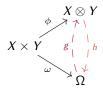
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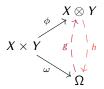


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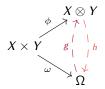
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- Then $g \circ h \circ \phi = \phi$, so $g \circ h = 1$ on $\phi(X \times Y)$, and hence on its span, $X \otimes Y$. Similarly, $h \circ g = 1$ on Ω, so $h^{-1} = g$.









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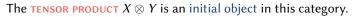
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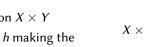
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The Definition of Lang

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Lang[3] defines the tensor product thusly:

Briefly wandering out of scope









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harkens to connection between universal arrows & universal objects in some comma category

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Tensor-Hom Adjunction



Notation

Write Hom(X, Y) for the SET OF LINEAR MAPS $X \to Y$. We abuse notation by interpreting Hom(X, Y) as a VECTOR SPACE when convenient.



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We can think of a bilinear map $\phi \in B(X \times Y, Z)$ instead as a linear map $\phi' \colon X \to Hom(Y, Z) \colon x \mapsto (y \mapsto \phi(x, y))$. In this way, we identify $B(X \times Y, Z)$ with Hom(X, Hom(Y, Z)).



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We can now succinctly formulate the linearizing property of $X \otimes Y$:

Theorem 7 (Tensor-hom adjunction) Hom $(X, \text{Hom}(Y, Z)) \cong \text{Hom}(X \otimes Y, Z)$. We say the tensor product $- \otimes Y$ is LEFT-ADJOINT to the internal hom functor Hom(Y, -).



Example (Cauchy stress tensor)

A common first-encounter with tensors is the CAUCHY STRESS TENSOR in classical mechanics. For direction vector $\mathbf{n} \in \mathbb{R}^3$, the 9-component stress tensor σ_{ij} gives the traction vector: $T_j^{(n)} = \sigma_{ij}n_i$.



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Example (Tensor product states)

In quantum mechanics, the set of states of a system is given by a Hilbert space. The two systems taken together have Hilbert space the tensor product of the individual systems' Hilbert spaces.

Entanglement arises precisely because not all tensors in this product are elementary.

Summary & Conclusion What we have shown



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This is succinctly captured by the **TENSOR-HOM ADJUNCTION**:

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We saw some examples of tensors, including from physics.

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